NASA Technical Memorandum 105837

11-3-92

Damping and Scattering of Electromagnetic Waves by Small Ferrite Spheres Suspended in an Insulator

Gerald W. Englert Lewis Research Center Cleveland, Ohio

August 1992



A

 $a^2\omega\mu/\tau$

DAMPING AND SCATTERING OF ELECTROMAGNETIC WAVES BY SMALL FERRITE SPHERES SUSPENDED IN AN INSULATOR

Gerald W. Englert
National Aeronautics and Space Administration
Lewis Research Center
Cleveland, Ohio 44135

SUMMARY

The intentional degradation of electromagnetic waves by their penetration into a media comprised of somewhat sparsely distributed energy absorbing ferrite spheres suspended in an electrical insulator is investigated. Results are presented in terms of generalized parameters involving wave length and sphere size, sphere resistivity, permeability, and spacing; and their influence on dissipation of wave power by eddy currents, magnetic hysteresis, and scattering is shown.

SYMBOLS

```
a^2 \omega \mu_o / \tau
A
       radius of sphere
       coefficients in equation (8)
\vec{\mathbf{B}}
       magnetic flux density
bo
       coefficients in equation (7)
       cosh (2a Re √jp)
       speed of wave travel
c_{o}
       \cos (2a \text{ Im } \sqrt{jp})
\overline{\mathbf{D}}
       electric displacement current
9
       constant defined in equation (6)
E
       electric field vector
       2.71828
Ħ
       magnetic inductance
H^{(1)}
       Hankel function of first kind
Im
       imaginary part of a complex quantity
J
       Bessel function of first kind
K
       power loss constants in equation (2)
k
       complex wave number
N
       number density of spheres per unit volume
P
       power
       parameter \omega \mu / \tau
p
Re
       real part of a complex quantity
```

```
S<sub>a</sub> sinh (2aRe √jp)
```

Shi hyperbolic sine integral

Si sine integral

s average spacing between center of spheres

 $s_a = \sin \left(2a \text{Im } \sqrt{jp}\right)$

t time

$$U_n - \mu_n - 1$$

$$U_{gn} |\mu_n|^2 + (1-2 \text{ Re } \mu_n)$$

x distance in direction of wave travel

$$lpha = 2 \; rac{{
m Re} \; \sqrt{{
m jp}}}{\sqrt{\mid {
m p} \mid}}$$

$$\beta = 2 \frac{\text{Im } \sqrt{\text{jp}}}{\sqrt{|\text{p}|}}$$

 λ wavelength

 μ permeability inside spheres, $\mu' - j\mu''$

 $\mu_{\rm n}$ normalized permeability, $\mu/\mu_{\rm o}$

 μ_{o} permeability of material in which spheres are immersed

au resistivity

 ω angular velocity

Subscripts

A applied wave

ec eddy current

hy hysteresis

Superscripts

' real part of complex quantity

" imaginary part of complex quantity

INTRODUCTION

Ferromagnetic oxides with their low electrical conductivity and large skin depth can be effective in absorbing electromagnetic wave energy (refs. 1 and 2). Their low weight densities, approximately 60 percent that of iron (ref. 3), make them attractive for flight applications. Spheres will be representative of the absorbent particles in the model compositions herein. Eddy currents are induced in the spheres by the applied electromagnetic waves resulting in dissipation of field energy by electrical resistance. Magnetic hysteresis phenomena cause additional losses. This evolves from the continual distortion of the crystalline structure of the Weiss domains and spin reorientation of the bound electrons induced by the waves (ref. 2, p. 16).

The intensity of interaction and resulting scattering of a plane electromagnetic wave by a conducting sphere falls off rapidly as the cube of the outside distance from the center of the sphere (ref. 4). This allows approximation of the multiple sphere problem by a succession of binary interactions except in the case of closely packed spheres.

Absorption of electromagnetic wave energy by small individual ferrite spheres was treated in reference 5. Two algebraic equations were derived expressing dimensionless power dissipation parameters for eddy currents and for hysteresis loss phenomena. These equations include complex magnetic permeability, μ , normalized by the permeability of the surrounding media, μ_0 , and a parameter of variables, $a^2\omega\mu_0/\tau$ involving angular wave frequency, ω , sphere radius, a, sphere resistivity, τ , as well as μ_0 . The present effort also includes power loss by inelastic scattering of the incident waves by the ferrite spheres and is treated by use of a complex wave number, k, (ref. 6).

Particles used in composites to isolate electronic equipment, as well as in wall coatings to damp unwanted electromagnetic radiation, are usually much smaller than the radiation wave length. At such low ratios of particle size to wave length diffraction plays a minor role and has been omitted.

ANALYSIS

The essential power balance equation for absorption of electromagnetic energy by means of eddy currents, P_{ec} , and magnetic hysteresis P_{hv} , induced in ferrite spheres can be expressed as

$$\frac{1}{2} \left[\overrightarrow{E}_{A} \cdot \frac{\partial \overrightarrow{D}_{A}}{\partial t} + \overrightarrow{H}_{A} \cdot \frac{\partial \overrightarrow{B}_{A}}{\partial t} \right] = \overrightarrow{H}_{A} \cdot \frac{\partial \overrightarrow{B}_{A}}{\partial t} = \frac{1}{2\mu_{A}} \frac{\partial \overrightarrow{B}_{A}^{2}}{\partial t} = -N(P_{ec} + P_{hy})$$
(3)

The applied electric and magnetic fields contribute equally in the initial power source. Here \vec{H}_A and \vec{B}_A are the magnetic induction and flux density of the wave whereas \vec{E}_A and \vec{D}_A are the electric field intensity and electric displacement vectors respectively. N is the number density of spheres per unit volume of composite. Power densities P_{ec} and P_{hy} are equal to constants K_{ec} and K_{hy} respectively times the square of the magnetic field of the applied wave, B_A , as in equations (33) and (35) of reference 5.

Including an additional power loss term, $P_s = B_A^2 K_s$ to represent elastic scattering of the incident wave, and therefore further loss of return signal, gives

$$\frac{dB_{A}^{2}}{B_{A}^{2}} = -\frac{2\mu_{A}N}{c_{o}} \left(K_{ec} + K_{hy} + K_{s} \right) dx$$
 (2)

where $c_0 = dx/dt$ is approximated by the speed of wave travel in the carrier material in which the spheres are immersed.

Integrating both sides of (2) gives the stopping power formula

$$B_{A}(x) = B_{A}(0)e^{-\mu_{A}N(K_{ec}+K_{hy}+K_{s})x/c}$$
(3)

By (ref. 5) the constant, Kec, for eddy currents, is

$$K_{ec} = \frac{P_{ec}}{B_{A}^{2}} = \frac{3\pi\omega^{2}a^{5}|\mu_{n}|^{2}|\mathscr{A}|}{\tau\mathscr{D}\operatorname{Re}\mathscr{A}} \left[S_{a} \operatorname{Im}\sqrt{j\mathscr{A}} + s_{a} \operatorname{Re}\sqrt{j\mathscr{A}} - \left(C_{a} - c_{a} \right) \frac{\operatorname{Re}\mathscr{A}}{|\mathscr{A}|} \right]$$
(4)

and for hysteresis

$$K_{hy} = \frac{P_{hy}}{B_{A}^{2}} = \frac{9\pi\omega^{2}a^{5}\mu''|\mu_{n}|\sqrt{|\mathcal{A}|}}{\tau\mathcal{D}} \left\{ \left(\frac{\alpha^{2} + \beta^{2}}{6} - 1 \right) \frac{C_{a} + c_{a}}{\sqrt{|\mathcal{A}|}} + \frac{1}{3} \left[\frac{\alpha S_{a} - \beta s_{a}}{|\mathcal{A}|} - \frac{C_{a} - c_{a}}{|\mathcal{A}|^{3/2}} + \frac{S_{a}}{\alpha} - \frac{s_{a}}{\beta} \right] + \left(\frac{4}{3} + \frac{\beta^{2}}{6} - \frac{\alpha^{2}}{2} \right) \alpha \operatorname{Shi}\left(2\operatorname{Re}\sqrt{\mathcal{A}}\right) - \left(\frac{4}{3} + \frac{\alpha^{2}}{6} - \frac{\beta^{2}}{2} \right) \beta \operatorname{Si}\left(2\operatorname{Im}\sqrt{\mathcal{A}}\right) \right\}$$

$$(5)$$

where parameter $\mathcal{A}=a^2\omega\mu/\tau$, normalized complex permeability $\mu_n=(\mu'-j\mu'')/\mu_o$, $j=\sqrt{-1}$, and μ_o is permeability of the media in which the spheres are suspended.

The constant \mathcal{D} in the denominator of equations (4) and (5) is

$$\mathcal{D} = U_{gn}^{2} |_{\mathscr{A}} | \left(C_{a} + c_{a} \right) + \left[U_{gn}^{2} + 2 \operatorname{Im} \left(\mathscr{A} \overline{\mu}_{n} \right) + |_{\mathscr{A}} |^{2} - 2 \operatorname{Im} \mathscr{A} \right] \left(C_{a} - c_{a} \right)$$

$$- \sqrt{2} \left\{ U_{gn}^{2} \left[\left(S_{a} + s_{a} \right) \operatorname{Re} \sqrt{\mathscr{A}} - \left(S_{a} - s_{a} \right) \operatorname{Im} \sqrt{\mathscr{A}} \right] \right]$$

$$- \left(S_{a} - s_{a} \right) \operatorname{Re} \left(U_{n} \mathscr{A} \sqrt{\mathscr{A}} \right) - \left(S_{a} + s_{a} \right) \operatorname{Im} \left(U_{n} \mathscr{A} \sqrt{\mathscr{A}} \right) \right\}$$

$$(6)$$

where $U_n = \mu_n - 1$, $U_{gn}^2 = |\mu_n|^2 + (1 - 2 \text{ Re } \mu_n)$, $p = \omega \mu / \tau$, $C_a = \cosh (2a \text{ Re } \sqrt{jp})$, $c_a = \cos (2a \text{ Im } \sqrt{jp})$, $S_a = \sinh (2a \text{ Re } \sqrt{jp})$, and $s_a = \sin (2a \text{ Im } \sqrt{jp})$.

Utilizing the scattering cross section, σ , of page 236 of reference 7 or page 449 of reference 8 gives

$$K_{s} = \frac{P_{s}}{B_{A}^{2}} = \frac{c\sigma}{2\mu_{A}} = \frac{\pi c}{\mu_{A}k^{2}} \sum_{\ell=1}^{\infty} (2\ell + 1)(|a_{\ell}|^{2} + |b_{\ell}|^{2})$$
 (7)

where a_{ℓ} and b_{ℓ} are expressed in Bessel functions of complex argument as

$$a_{\ell} = -\frac{J_{\ell+1/2}(ka)}{H_{\ell+1/2}^{(1)}(ka)}, \quad b_{\ell} = -\frac{(\ell-1)J_{\ell+1/2}(ka) - kaJ_{\ell-1/2}(ka)}{(\ell-1)H_{\ell+1/2}^{(1)}(ka) - kaH_{\ell-1/2}^{(1)}(ka)},$$
(8)

and

$$k^{2} = \left(\frac{\omega}{c}\right)^{2} + \frac{j\omega\mu}{\tau} = \left(\frac{2\pi}{\lambda}\right)^{2} + \frac{\omega}{\tau}(\mu'' + j\mu').$$

Three coefficients were used in each of the rapidly converging series of equation (7). They can be expressed as

$$a_1 = -\frac{\sin(ka) - ka \cos(ka)}{ka[\sin(ka) - j \cos(ka)]}, \qquad (9)$$

$$a_{2} = \frac{\frac{3}{ka} \cos(ka) + \left[1 - \frac{3}{(ka)^{2}}\right] \sin(ka)}{\frac{3}{ka} \cos(ka) + \left[1 - \frac{3}{(ka)^{2}}\right] \sin(ka) + j \left\{\frac{3}{ka} \sin(ka) - \left[1 - \frac{3}{(ka)^{2}}\right] \cos(ka)\right\}},$$
 (10)

$$a_{3} = \frac{\left[1 - \frac{5!}{8(ka)^{2}}\right] \cos(ka) + \left[\frac{5!}{8(ka)^{3}} - \frac{3!}{ka}\right] \sin(ka)}{\left[1 - \frac{5!}{8(ka)^{2}} + j\left(\frac{3!}{ka} - \frac{5!}{8(ka)^{3}}\right)\right] \left[\cos(ka) + j \sin(ka)\right]},$$
(11)

$$b_{1} = \frac{\left[1 - (ka)^{2}\right] \sin(ka) - ka \cos(ka)}{e^{jka} \left\{ka - j\left[(ka)^{2} - 1\right]\right\}}$$
(12)

$$b_{2} = \frac{\left[3(ka)^{2} - 6\right]\sin(ka) - \left[(ka)^{3} - 6ka\right]\cos(ka)}{e^{jka}\left\{(ka)^{3} + 6ka + j\left[6 - (ka)^{2}\right]\right\}},$$
(13)

and

$$b_{3} = \frac{(ka)^{4} \sin(ka) + 6(ka)^{3} \cos(ka) - 21(ka)^{2} \sin(ka) - 45[ka \cos(ka) - \sin(ka)]}{e^{jka} \left\{ \frac{9}{2} (ka)^{3} + 45ka - j \left[(ka)^{4} - 18(ka)^{3} - 3(ka)^{2} + 45ka \right] \right\}}$$
(14)

by use of pages 966 and 967 of reference 9.

Next consider the generalized distance to reduce $B_A(x)/B_A(0)$ to 1/e and let number density $N=1/s^3$, where s is the average spacing between centers of spheres. The exponent in equation (3) written in terms of dimensionless parameters, then reduces to -1 when

$$\left(\frac{a}{s}\right)^{3} \frac{\omega x}{c} = 1 / \left\{ \frac{\tau \left(P_{ec} + P_{hy}\right)}{\omega^{2} a^{5} B^{2}} \left(a^{2} \omega \mu_{A} / \tau\right) + \frac{\lambda / a}{2(ka)^{2}} \sum_{1}^{3} (2\ell + 1) \left(|a_{\ell}|^{2} + |b_{\ell}|^{2}\right) \right\}$$
(15)

The individual contributions of eddy currents and hysteresis are given in reference 5 over wide ranges of \mathcal{A}_{o} .

RESULTS

The generalized distance of wave travel, $(a/s)^3 \omega x/c$, in the composite to reduce $B_A(x)/B_A(0)$ to 1/e is plotted versus μ'/μ_o and μ''/μ_o for fixed values of $a^2\omega\mu_o/\tau$ on three-dimensional plots of figures 1(a) to 2(h). The range of generalized distance decreases with increase of parameter $a^2\omega\mu_o/\tau$ for absorption (figs. 1(a) to (d)) as well as for scattering when holding λ/a constant: figures 2(a), (e), and (g) with $\lambda/a=10^3$ and figures 2(b), (f), and (h) with $\lambda/a=10^5$.

The peaks of the surfaces move to higher μ'/μ_0 as $a^2\omega\mu_0/\tau$ and/or λ/a decrease. Pronounced concave trough-like contours indicate the minimum distance of absorption whereas evident convex patterns give the maximum distances for reduction of scattering. The dashed lines shown are merely diagonals where $\mu'=\mu''$.

As λ/a is decreased from 10^5 the ridges gradually rotate their directions from diagonal to ones more closely parallel to the μ''/μ_0 axis. With further decrease of λ/a the contours flatten out to nearly constant values of generalized distance. This is shown over an especially wide range of λ/a in figures 2(c) to (f). This trend also appears as $a^2\omega\mu_0/\tau$ is reduced with λ/a held constant. For example, note the change in curve pattern between figures 2(g), (e), and (a).

The generalized distance to reduce $B_A(x)/B_A(0)$ to 1/e by absorption is nonsymmetric about the dashed line due to the lack of symmetry in the contribution from magnetic hysteresis. The results for scattering are symmetric when like values are selected for μ' and μ'' .

CONCLUDING REMARKS

By use of generalized parameters the seven independent variables: sphere radius a, complex permeability μ' -j μ'' , resistivity τ , mean distance between spheres s, electromagnetic field strength B_A , and wave length λ times frequency combine to (1) a generalized distance of wave travel, $(a/s)^3 \omega x/c$, to reduce the applied magnetic field ratio, $B_A(x)/B_A(0)$ to 1/e of its initial strength, (2) a power absorption parameter $a^2\omega\mu_o/\tau$, and (3) magnetic permeability ratios μ'/μ_o and μ''/μ_o .

The range of generalized distance for reducing $B_A(x)/B_A(0)$ to 1/e by <u>absorption</u> decreases with increase of $a^2\omega\mu_o/\tau$ over its range of 10^{-7} to 10 while the peaks of these plots reduce from the order of 10^7 to 10^3 . The dashed lines shown are merely diagonals where $\mu'=\mu''$.

The three-dimensional plots of generalized distance by <u>scattering</u>, however, are symmetrical about the dashed lines. Also an additional parameter, wave length divided by sphere radius, λ/a , enters as an independent variable for these plots.

REFERENCES

- 1. VonAulock, W.H.: Handbook of Microwave Ferrite Materials. New York: Academic Press, 1965.
- 2. Smit, J.; and Wijn, H.P.J.: Ferrites. New York: John Wiley & Sons, 1959.
- Schlicke, H.M.: Essentials of Dielectromagnetic Engineering. New York: John Wiley & Sons, 1961, p. 75.
- 4. Smythe, W.R.: Static and Dynamic Electricity. 3rd ed., New York: McGraw-Hill, 1968, p. 377.
- Englert, G.W.: Parametric Study of Power Absorption from Electromagnetic Waves by Small Ferrite Spheres. NASA TP-2949, 1989.
- Stratton, J.D.: Electromagnetic Theory. New York: McGraw-Hill, 1941, p. 392.
- 7. Panofsky, W.K.H.; and Phillips, M.: Classical Electricity and Magnetism. 2nd ed., Cambridge, MA: Addison Wesley, 1955, p. 236.
- 8. Kerr, D.E.: The Radar Cross Section of Isolated Targets. Propagation of Short Radio Waves, D.E. Kerr, ed., M.I.T. Radiation Laboratory Series 13, 1951, p. 449.
- 9. Gradshteyn, I.S.; and Ryzhik, I.M.: Table of Integrals, Series, and Products. New York: Academic Press, 1965, pp. 966-967.

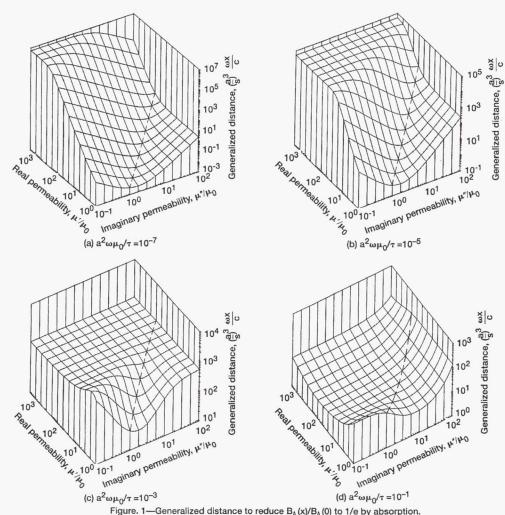
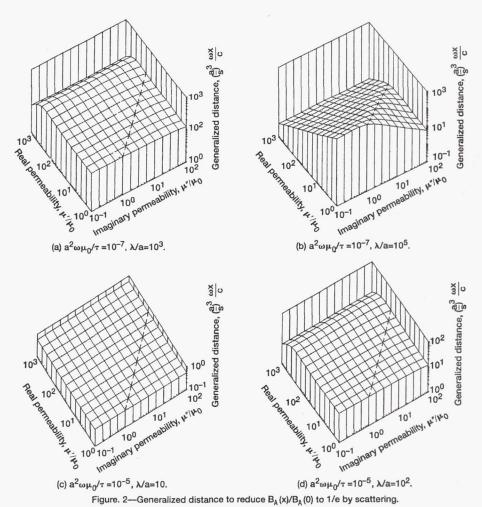
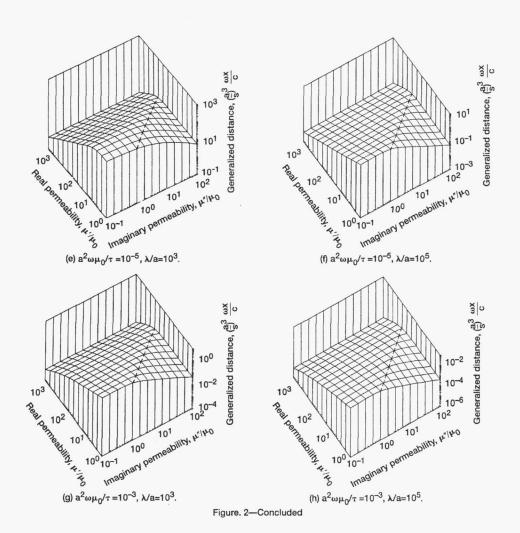


Figure. 1—Generalized distance to reduce $B_{\underline{A}}\left(x\right)\!/B_{\underline{A}}\left(0\right)$ to 1/e by absorption.





REPORT DOCUMENTATION PAGE

Form Approved
OMB No. 0704-0188

Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing the collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden, to Washington Headquarters Services, Directorate for information Operations and Reports, 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302, and to the Office of Management and Budget, Paperwork Reduction Project (0704-0188), Washington, DC 20503.

1. AGENCY USE ONLY (Leave blank)	The state of the s	The Control of the Co	
	August 1992		chnical Memorandum
4. TITLE AND SUBTITLE			5. FUNDING NUMBERS
Damping and Scattering of Electromagnetic Waves by Small Ferrite			
Spheres Suspended in an Insulator			
6. AUTHOR(S)			WU-505-62-52
Gerald W. Englert			
Gerald W. Eligiert			
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES)			8. PERFORMING ORGANIZATION
			REPORT NUMBER
National Aeronautics and Space Administration			
Lewis Research Center			E-6363-1
Cleveland, Ohio 44135–3191			
O CROMCODING MONITORING ACENCY MANEC(C) AND ADDRESS(ES)			40 ORGANGORING/MONITORING
9. SPONSORING/MONITORING AGENCY NAMES(S) AND ADDRESS(ES)			10. SPONSORING/MONITORING AGENCY REPORT NUMBER
National Aeronautics and Space Administration			
Washington, D.C. 20546–0001			NASA TM-105837
			111111111111111111111111111111111111111
11. SUPPLEMENTARY NOTES			
Gerald W. Englert, NASA Lewis Research Center, Cleveland, Ohio. Responsible person, Gerald W. Englert,			
(216) 433–5887.			
12a. DISTRIBUTION/AVAILABILITY S	TATEMENT	Ţ.	12b. DISTRIBUTION CODE
Table Storms Trought Alexander Transaction			ZB. BIOTHIBOTION COSE
Unclassified - Unlimited			
Subject Category 32			
13. ABSTRACT (Maximum 200 words)			
The intentional degradation of electromagnetic waves by their penetration into a media comprised of somewhat			
sparsely distributed energy absorbing ferrite spheres suspended in an electrical insulator is investigated. Results are presented in terms of generalized parameters involving wave length and sphere size, sphere resistivity, permeability,			
and spacing; and their influence on dissipation of wave power by eddy currents, magnetic hysteresis and scattering			
is shown.			
,			
14. SUBJECT TERMS			15. NUMBER OF PAGES
Magnetics			12
5			16. PRICE CODE
,		7	A03
17. SECURITY CLASSIFICATION OF REPORT	18. SECURITY CLASSIFICATION OF THIS PAGE	19. SECURITY CLASSIFICAT OF ABSTRACT	TION 20. LIMITATION OF ABSTRACT
Unclassified	Unclassified	Unclassified	